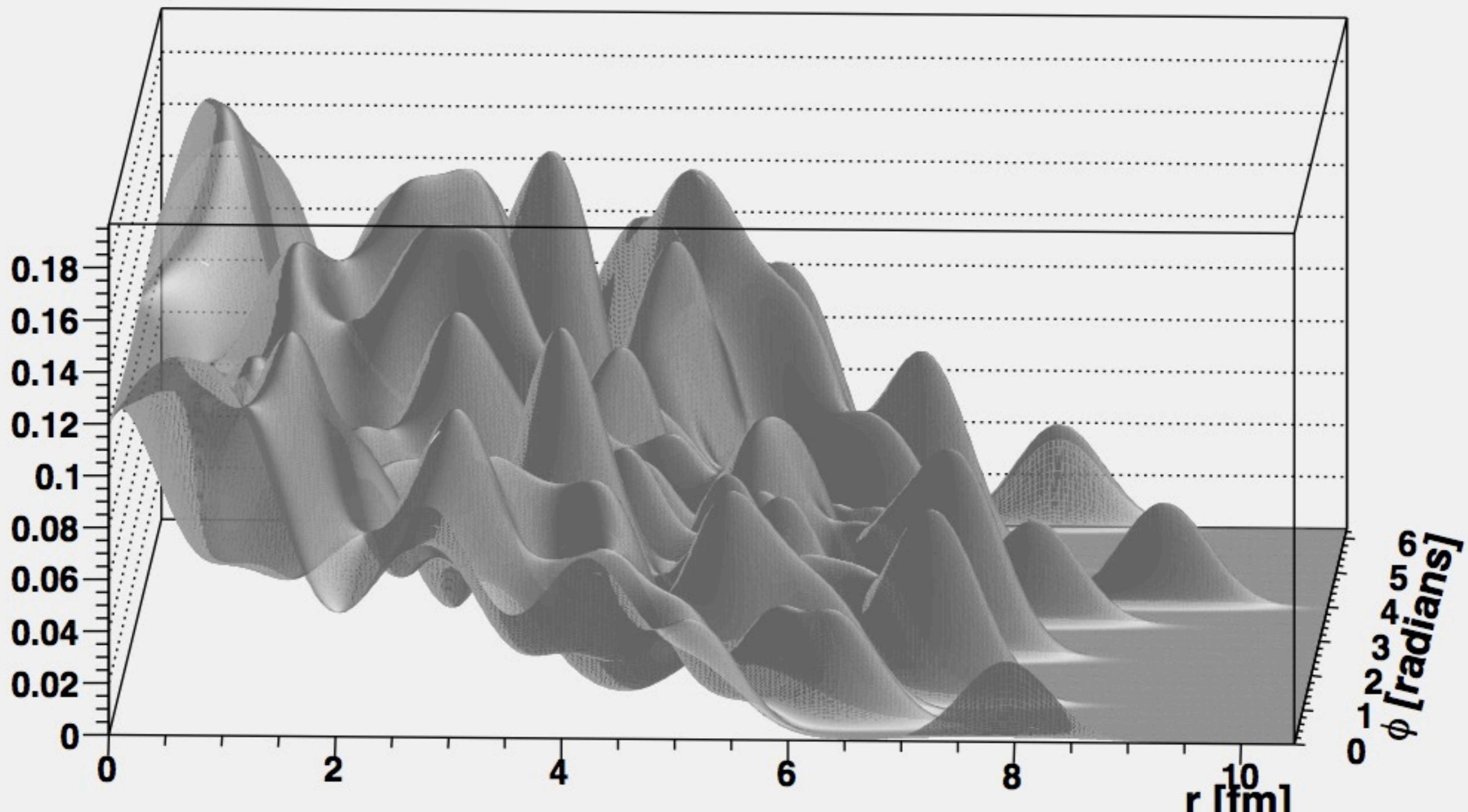
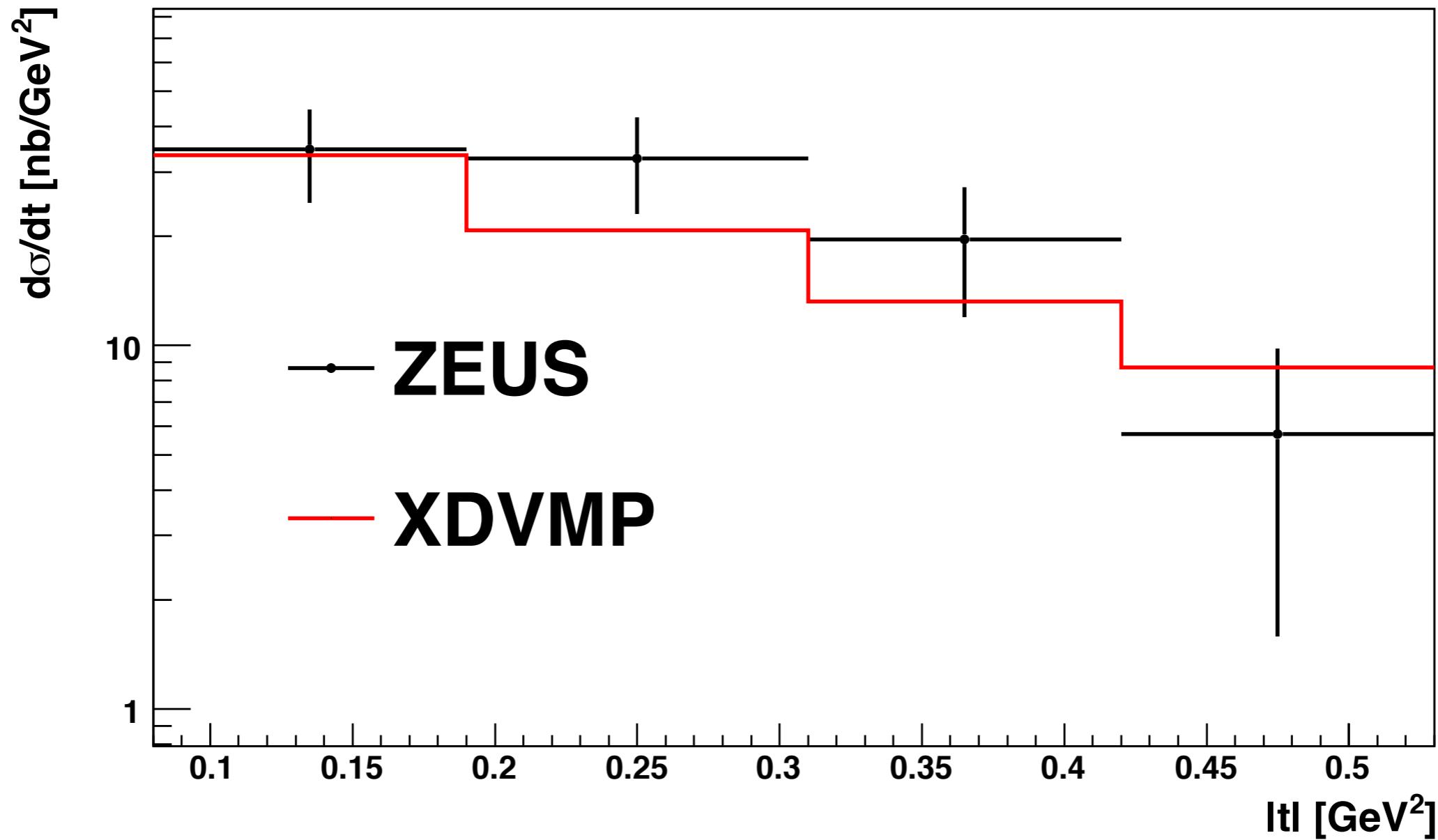


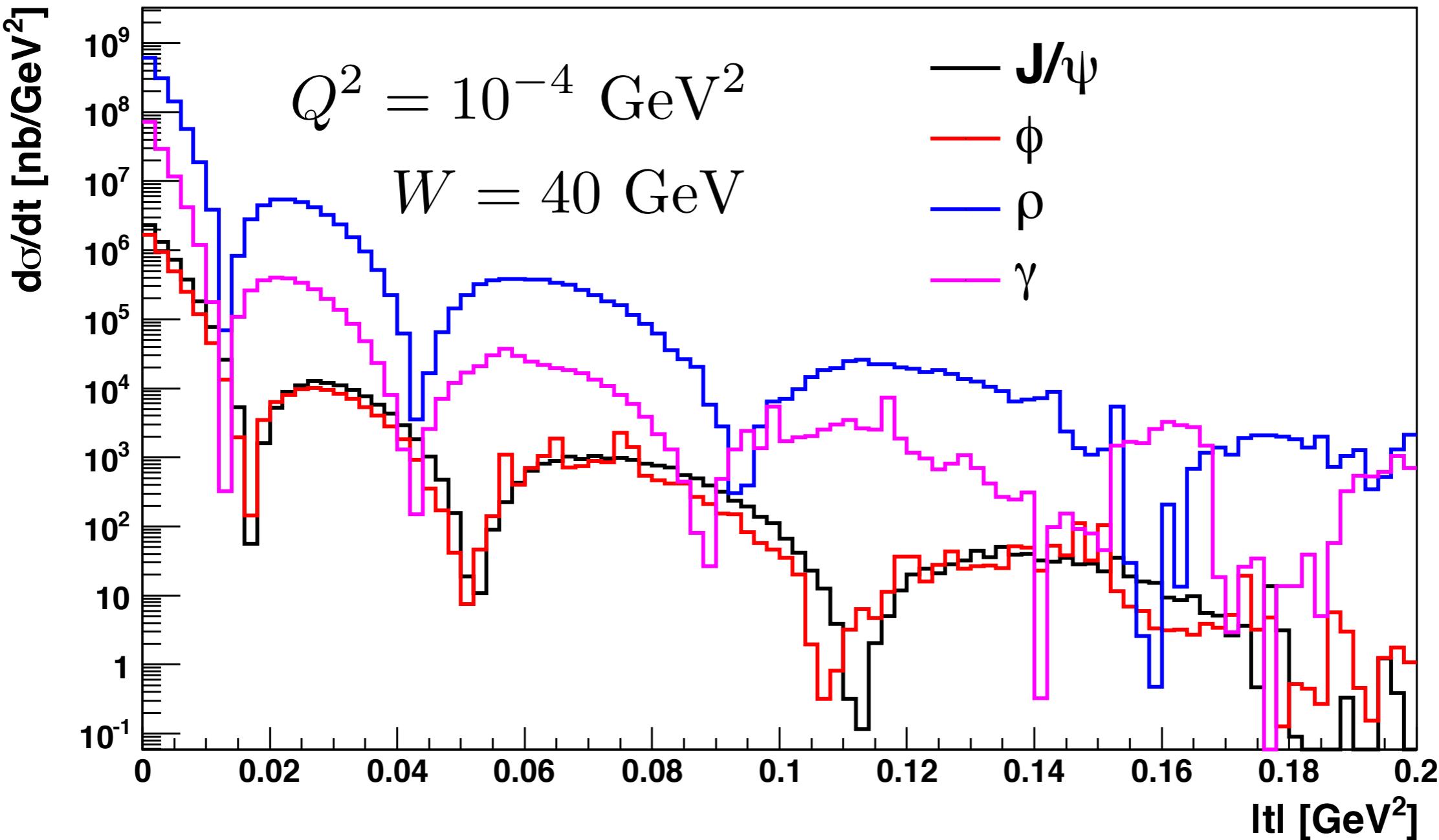
Update on XDVMP for eA



Tobias Toll
02/17/11

DVCS in ep bCGC

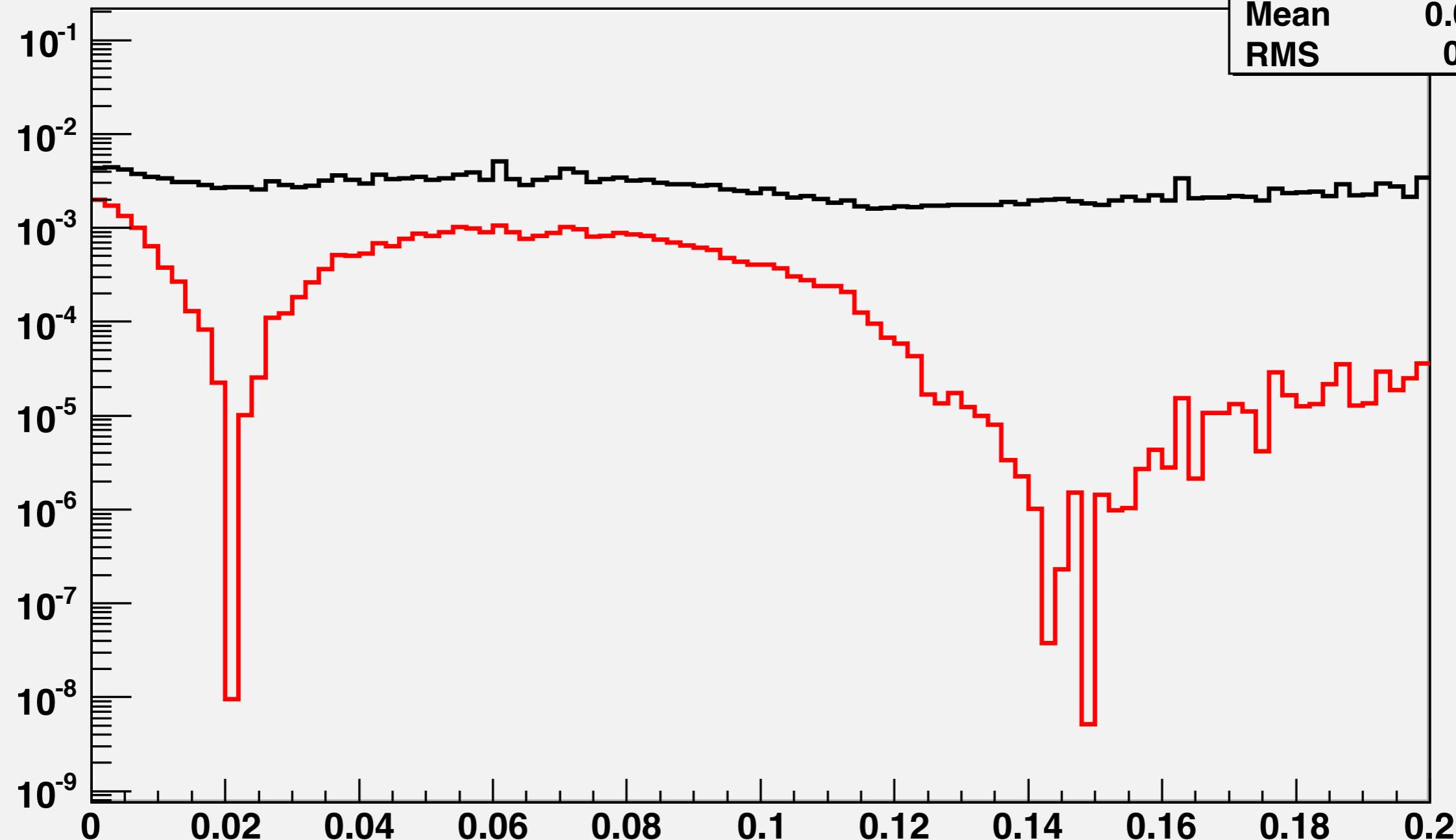


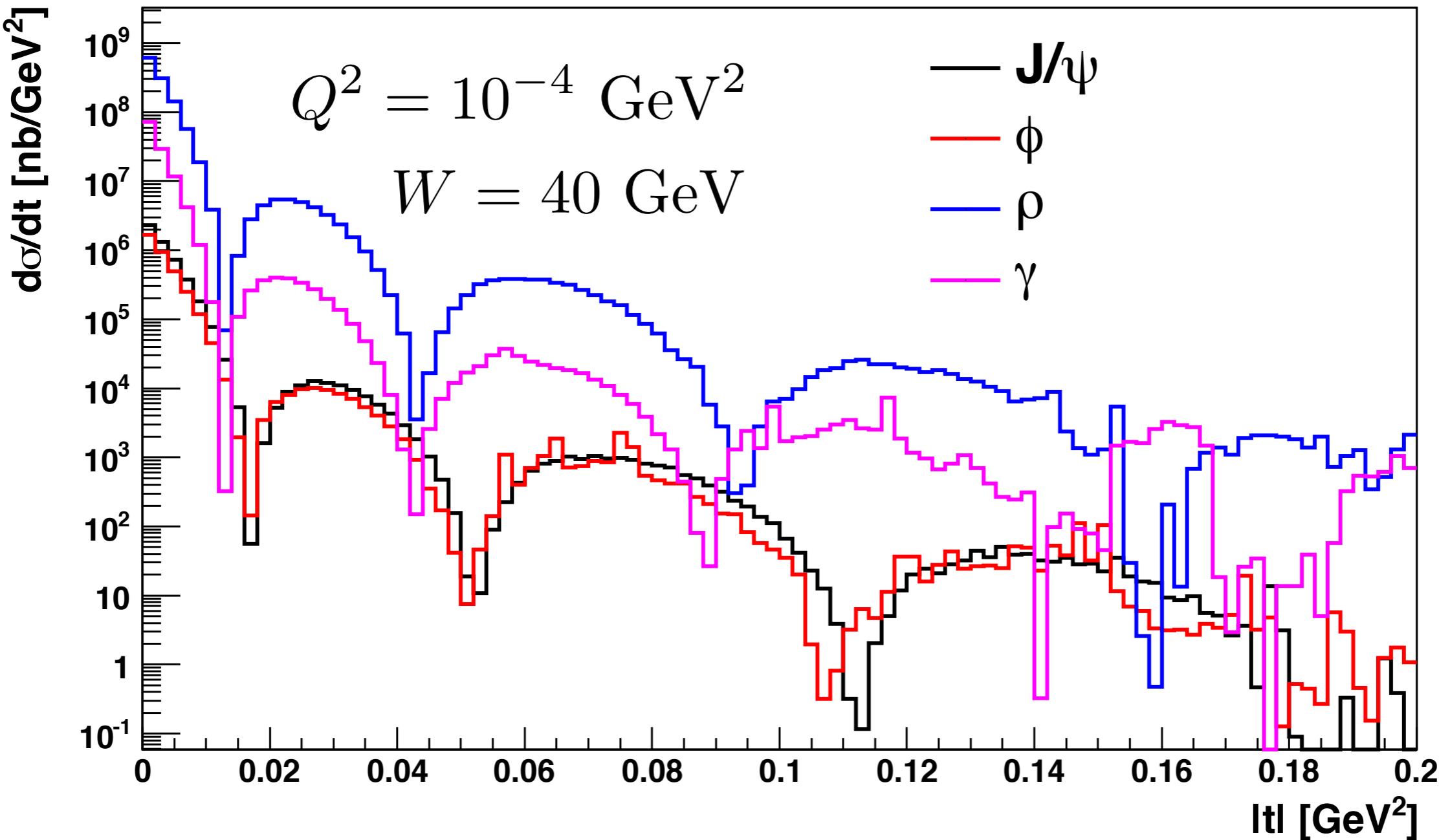


dsigmadt

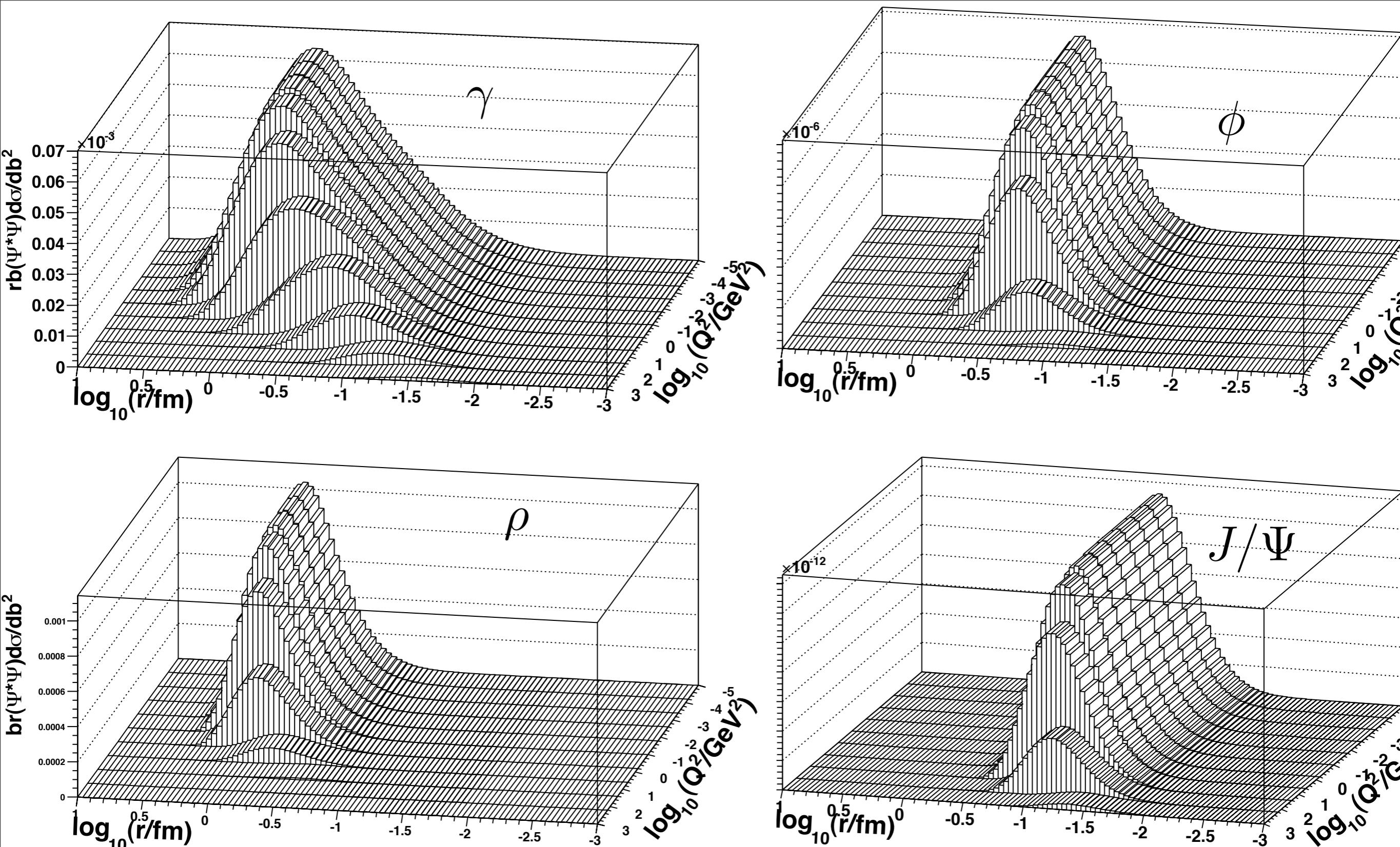
$W = 40 \text{ GeV}$ $Q^2 = 500 \text{ GeV}^2$

coherent	
Entries	100
Mean	0.05694
RMS	0.0346





$$\begin{aligned}
 \mathcal{A} = & \int dr \int dz \int db \int d\phi \\
 & \frac{1}{2} rb (\Psi_V^* \Psi) (r, z) \frac{d^2 \sigma_{q\bar{q}}}{db d\phi} (x, r, b, \phi, \Omega) J_0([1-z]r\Delta) \cos(b\Delta \cos\phi)
 \end{aligned}$$



$$\frac{1}{2} rb(\Psi_V^* \Psi)(r, z) \frac{d^2 \sigma_{q\bar{q}}}{db d\phi}(x, r, b, \phi, \Omega) J_0([1-z]r\Delta) \cos(b\Delta \cos\phi)$$

Premature Discussion

For measuring the spacial gluon distribution in nuclei, it's important to minimize the contributions from the term $J_0([1 - z]r\Delta) \text{ vs. } \cos(b\Delta \cos \phi)$

This seems to be possible in kinematic regimes dominated by large r , making the first term oscillate faster than the second term:

DVCS in photoproduction not J/Psi

This is contrary to what Caldwell/Kowalski says in their paper... and may contradict Thomas' results with the inverse Fourier transform so: I may very well be wrong

hence: **More Investigations needed!!**